# STABILITY OF THE TRIVIAL SOLUTION OF A SYSTEM OF TWO LINEAR DIFFERENTIAL EQUATIONS WITH PERIODIC COEFFICIENTS 

## (OB USTOICHIVOSTI TBIVIAL' NOGO RESBENIIA SISTEUY DVUEH LINEINYIH DIFPERENTSIAL'NYEH URAVNENII S PERIODICHESEIMI KOEFFITSIENTAMI)

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We propose to find a new estimate of the characteristic constant of Liapunov [1] for the system mentioned in the title when the functions on the secondary diagonal of the coefficient matrix are of definite and the same sign.

1. Let the equations of the first approximation of the perturbed motion of a dynamical system be given in the form

$$
\begin{equation*}
\frac{d x_{1}}{d t}=a_{11}(t) x_{1}+a_{12}(t) x_{2}, \quad \frac{d x_{2}}{d t}=a_{21}(t) x_{1}+a_{22}(t) x_{2} \tag{1}
\end{equation*}
$$

Where $a_{i j}(t)(i, j=1,2)$ are real, integrable, sectionally continuous periodic functions of period $\omega$. We introduce the notation

$$
\int_{0}^{\infty} a_{11}(t) d t=\alpha, \quad \int_{0}^{\infty} a_{22}(t) d t=\beta
$$

and, without affecting the generality, we assume that $a>\beta$. If $a+\beta>0$, the unperturbed motion of the dynamical system for which the equation of the first approximation of the perturbed motion has the form (1), is stable by a well-known formula of Liapunov. In the sequel we shall assume that $a+\beta<0$, which, in view of the fact that $a-\beta>0$, can be written in the form

$$
\begin{equation*}
\beta \leqslant \alpha \leqslant-\beta \quad(\beta \leqslant 0) \tag{2}
\end{equation*}
$$

Making the change of variables

$$
x_{1}=x \exp \int_{0}^{t} a_{11}\left(t_{1}\right) d t_{1}, \quad x_{2}=y \exp \int_{0}^{t}\left[a_{22}\left(t_{1}\right)+\frac{\alpha-\beta}{\omega}\right] d t_{1}
$$

we transform the system (1) to the form

$$
\begin{equation*}
\frac{d x}{d t}=r(t) y, \quad \frac{d y}{d t}=q(t) x-a y \quad\left(a=\frac{a-\beta}{\omega} \geqslant 0\right) \tag{3}
\end{equation*}
$$

where

$$
\begin{aligned}
& r(t)=a_{12}(t) \exp \int_{0}^{t}\left[a_{22}\left(t_{1}\right)-a_{11}\left(t_{11}+a\right] d t_{1}\right. \\
& q(t)=a_{21}(t) \exp \int_{0}^{t}\left[a_{11}\left(t_{1}\right)-a_{22}\left(t_{1}\right)-a\right] d t_{1}
\end{aligned}
$$

The characteristic equation for the system (3) has the form

$$
p^{2}-2 A p+e^{-a \omega}=0 \quad\left(A=\frac{1}{2} \operatorname{sp} X(\omega) \quad X(0)=I_{2}\right)
$$

where $X(t)$ is the matrix of the system (3). The characteristic constant $A$, in the sense of Liapunov [1], can be represented by a converging power series

$$
\begin{equation*}
A=A_{0}+A_{1}+A_{2}+\cdots \tag{4}
\end{equation*}
$$

The formulas for the coefficients $A_{n}$ were derived* in [2].

$$
\begin{gather*}
A_{0}=\frac{1}{2}\left(1+e^{\beta-\alpha}\right), \quad A_{1}=\frac{1}{2} \int_{0}^{\omega} e^{a t_{1}}\left[\psi(\omega)-\left(1-e^{-a \omega}\right) \psi_{1}\right] q_{1} d t_{1} \\
A_{n}=\frac{1}{2} \int_{0}^{\omega} d t_{1} \int_{0}^{t_{1}} d t_{2} \ldots \int_{0}^{t_{n-1}} e^{a t_{1}}\left(\psi_{1}-\psi_{2}\right) \ldots  \tag{5}\\
\ldots e^{a t_{n-1}}\left(\psi_{n-1}-\psi_{n}\right) e^{a t_{n}}\left[\psi(\omega)-\psi_{1}+\psi_{n} e^{-a \omega}\right] q_{1} \ldots q_{n} d t_{n}
\end{gather*}
$$

where

$$
\psi(t)=\int_{0}^{t} e^{-a t_{1 r}}\left(t_{1}\right) d t_{1}, \quad \psi_{i}=\psi\left(t_{i}\right), \quad q_{i}=q\left(t_{i}\right)
$$

By means of an orthogonal transformation with periodic coefficients (see [3]) one can reduce the system (1) to a form in which the functions $a_{12}$ and $a_{21}$ (and hence the functions $r$ and $q$ ) are of a definite sign.

In [4], the functions $r$ and $q$ were assumed to be of definite but opposite signs.
2. From here on we shall always assume that $r$ and $q$ are of definite sign on $[0, \omega]$ (furthermore, $r(t) q(t) \neq 0$ for $0 \leqslant t<\omega$, because otherwise the system (3) can be integrated over intervals).

- In [2] the constant $A$ stood for $\operatorname{sp} X(\omega) /\left(1+e^{-\alpha c y}\right.$.

In this case it follows from (5) that $A_{\nu}>0(\nu=1,2, \ldots)$ and the series (4) is positive-definite.

In [2] it was shown that if $a+\beta \leqslant 0$, the unperturbed motion of the dynamical systemk for which the equation of the first approximation has the form (1), is unstable if $|A|>\left(e^{-\alpha}+e^{\beta}\right) / 2$. Since $A>A_{0}$ here, it is sufficient to require for instability that $A_{0}>\left(e^{-a}+e^{\beta^{0}}\right) / 2$, which in view of the first formula of (5) can be written in the form

$$
\left(1-e^{-\alpha}\right)\left(1-e^{\beta}\right) \geqslant 0
$$

Recalling (2) we conclude that if $a>0$ then the last inequality is satisfied. Hence, we have proved the following theorem.

Theorem. Suppose that in the equations of the first approximation of the perturbed motion (1) the functions $a_{12}$ and $a_{21}$ are of definite sign (i.e. retaining the same sign, they may become zero but in such a way that their product does not vanish identically on ( $0, \omega$ ) and is of one sign). If, furthermore, the mean value of one of the functions $a_{11}$ or $a_{22}$ is non-negative, then the unperturbed motion is unstable.

Consequence. If in the equation

$$
\begin{equation*}
\ddot{x}+s(t) \ddot{x}+p(t) x=0, \quad(s(t+\omega)=s(t), \quad p(t+\omega)=p(t)) \tag{6}
\end{equation*}
$$

the periodic coefficient $p(t)$ is non-positive for all values of $t(0 \leqslant t \leqslant \omega)$, then the trivial solution of Equation (6) is unstable irrespective of the behavior of the second periodic coefficient $s(t)$.

For the proof it is sufficient to write Equation (6) in the form of the system

$$
\frac{d x}{d t}=\dot{x} \quad \frac{d \dot{x}}{d t}=-p(t) x-s(t) \dot{x}
$$

and to apply the established theorem.
3. The case that remains to be considered is the one when $\beta \leqslant a<0$. In this case the undisturbed motion will be stable if

$$
\begin{equation*}
A<\frac{1}{2}\left(e^{-\alpha}+e^{\beta}\right) \tag{7}
\end{equation*}
$$

and it will be unstable if the inequality sign is reversed [2].
In [2] there were proved certain inequalities which can be written for the terms of the series (4) in the following way:
4. Because of the first of the inequalities (8) with $n=2,3, \ldots$,
we have that

$$
A_{2}<A_{1}^{2}, \quad A_{3}<\frac{2}{3} A_{1} A_{2}<\frac{2}{3} A_{1}^{3}, \ldots, A_{n}<\frac{2^{n-1}}{n!} A_{1}{ }^{n}
$$

Hence, we find that

$$
A<A_{0}+A_{1}+\frac{1}{2} \sum_{n=2}^{\infty} \frac{2^{n}}{n!} A_{1}^{n}=\frac{1}{2}\left[e^{\beta-\alpha}+\frac{1}{2} e^{2 A_{2}}\right.
$$

Under the condition that

$$
\frac{1}{2} e^{\beta-\alpha}+\frac{1}{2} e^{2 A_{1}} \leqslant \frac{1}{2}\left(e^{-\alpha}+e^{\beta}\right) \text { или } A_{1} \leqslant \frac{1}{2} \ln \left(e^{-\alpha}+e^{\beta}-e^{\beta-\alpha}\right)
$$

the inequality (i: will hold, and, hence, the unperturbed motion will be stable.*
5. Let us now apply to the above estimates the second inequality of (8)
$A_{2}<\frac{1}{4} e^{\alpha-\beta} A_{1}{ }^{2}, \quad A_{3}<\frac{1}{9} e^{\alpha-\beta} A_{1} A_{2}<\frac{1}{36} e^{2(\alpha-\beta)} A_{1}{ }^{3} \ldots, A_{n}<\frac{1}{(n!)^{2}} e^{(n-1)(\alpha-\beta)} A_{1}{ }^{n}$
We thus obtain
$A<A_{0}+A_{1}+\sum_{n=2}^{\infty} \frac{1}{(n!)^{2}} e^{(n-1)(\alpha-\beta)} A_{1}^{n}=\frac{1}{2}\left(1-e^{\beta-\alpha}\right)+e^{\beta-\alpha} I_{0}\left(2 \sqrt{e^{\alpha-\beta} A_{1}}\right)$
where $I_{0}(z)$ is a Bessel function of an imaginary argument. If

$$
\frac{1}{2}\left(1-e^{\beta-\alpha}\right)+e^{\beta-\alpha} I_{0}\left(2 \sqrt{e^{\alpha-\beta} A_{1}}\right) \leqslant \frac{1}{2}\left(e^{-\alpha}+e^{\beta}\right)
$$

i.e.

$$
\begin{equation*}
I_{0}\left(2 \sqrt{e^{\alpha-\beta} A_{1}}\right) \leqslant \frac{1}{2}\left(1+e^{\alpha}+e^{-\beta}-e^{\alpha-\beta}\right) \tag{10}
\end{equation*}
$$

then the unperturbed motion is stable.** The inequality (10) is preferable over (9) when $\alpha=\beta<\ln 4$.

* In the last formula on p. 136 of [2] the denominator of the logarithmic term should be dropped.
** In Equation (11.5) of [2] the right-hand side should be written in the form

$$
\frac{1}{2}\left(1+e^{\alpha}+e^{-\beta}-e^{\alpha-\beta}\right)
$$

and the index of the exponent of the left-hand side in the form $a-\beta$.

If we begin the estimates of the terms of the series (4) with $A_{3}$, we find that if either one of the following inequalities

$$
\begin{equation*}
A_{2} \leqslant \frac{1}{e^{2 A_{1}}-1-2 A_{1}}\left[\left(e^{-\alpha}-1\right)\left(1-e^{\beta}\right)-2 A_{1}\right] A_{1}^{2} \tag{11}
\end{equation*}
$$

or

$$
\begin{equation*}
A_{2} \leqslant \frac{\left[\left(e^{-\alpha}-1\right)\left(1-e^{\beta}\right)-2 A_{1}\right] e^{2(\alpha-\beta)} A_{1}{ }^{2}}{8\left[I_{0}\left(2 \sqrt{e^{\alpha-\beta} A_{1}}\right)-1-e^{\alpha-\beta} A_{1}\right]} \tag{12}
\end{equation*}
$$

holds then the undisturbed motion is stable.
6. In [4] there was proved an inequality which for the terms of the series (4) can be written in the form

$$
\begin{equation*}
\frac{A_{n+1}}{A_{n}} \leqslant \frac{n}{n+1} \frac{A_{n}}{A_{n-1}} \quad(n \geqslant 1) \tag{13}
\end{equation*}
$$

If we set $n=1,2,3$ we obtain, respectively

$$
\frac{A_{2}}{A_{1}} \leqslant \frac{1}{2} \frac{A_{1}}{A_{0}}, \quad \frac{A_{3}}{A_{2}} \leqslant \frac{2}{3} \frac{A_{2}}{\overline{A_{1}}}, \quad \frac{A_{4}}{A_{3}} \leqslant \frac{3}{4} \frac{A_{3}}{A_{2}}, \ldots
$$

Making use of the products of the last inequalities, we obtain the result for the sum of the series (4):

$$
\begin{gather*}
A \leqslant A_{0}+A_{1}+\frac{1}{2} \frac{A_{1}^{2}}{A_{0}}+\frac{1}{3} \frac{A_{1} A_{2}}{A_{0}}+\frac{1}{4} \frac{A_{1} A_{3}}{A_{0}}+\ldots= \\
=A_{0}+\frac{2}{3} A_{1}+\frac{1}{6} \frac{A_{1}^{2}}{A_{0}}+\frac{A_{1}}{A_{0}}\left[\frac{1}{3} A_{0}+\frac{1}{3} A_{1}+\frac{1}{3} A_{2}+\frac{1}{4} A_{3}+\ldots\right] \tag{14}
\end{gather*}
$$

The sum of the series within the brackets is less than $A / 3$, and hence we have

$$
A\left(1-\frac{1}{3} \frac{A_{1}}{A_{0}}\right)<A_{0}+\frac{2}{3} A_{1}+\frac{1}{6} \frac{A_{1}{ }^{2}}{A_{0}}
$$

If $A_{1}>3 A_{0}$, then the last inequality is obviously valid. If, however, $A_{1}<3 A_{0}$, then
and under the condition

$$
A<\frac{6 A_{0}^{2}+4 A_{1} A_{0}+A_{1}^{2}}{6 A_{0}-\frac{A_{1}}{2 A_{1}}}
$$

$$
\begin{equation*}
\frac{6 A_{0}^{2}+4 A_{1} A_{0}+A_{1}^{2}}{3 A_{0}-A_{1}} \leqslant e^{-\alpha}+e^{\beta} \tag{15}
\end{equation*}
$$

the unperturbed motion is stable.
Let us now select from the right-hand side of the inequality (14) all terms up to and including $A_{2}$

$$
A \leqslant A_{0}+\frac{3}{4} A_{1}+\frac{1}{4} \frac{A_{1}^{2}}{A_{0}}+\frac{1}{12} \frac{A_{1} A_{2}}{A_{0}}+\frac{A_{1}}{A_{0}}\left[\frac{1}{4} A_{0}+\frac{1}{4} A_{1}+\frac{1}{4} A_{2}+\frac{1}{4} A_{3}+\ldots\right]
$$

We thus obtain

$$
\begin{equation*}
A\left(1-\frac{1}{4} \frac{A_{1}}{A_{0}}\right)<A_{0}+\frac{3}{4} A_{1}+\frac{1}{4} \frac{A_{1}^{2}}{A_{0}}+\frac{1}{12} \frac{A_{1} A_{2}}{A_{0}} \tag{16}
\end{equation*}
$$

which, obviously, will be true if $A_{1}>4 A_{0}$. If $A_{1}<4 A_{0}$, then, under the condition that

$$
\frac{12 A_{0}^{2}+9 A_{1} A_{0}+3 A_{1}^{2}+A_{1} A_{2}}{4 A_{0}-A_{1}}<\frac{3}{2}\left(e^{-\alpha}+e^{\beta}\right)
$$

the unperturbed motion is stable.
7. In regard to the sufficient conditions for instability, it can be said that they can be found more simply. Indeed, $A>A_{0}+A_{1}$ and, therefore, if $\left.A_{0}+A_{1}>e^{-a}+e^{\beta}\right) / 2$ we have also that $A>\left(e^{-a}+e^{\beta}\right) / 2$. Hence, if the inequality

$$
\begin{equation*}
A_{1} \geqslant \frac{1}{2}\left(e^{-\alpha}-1\right)\left(1-e^{\beta}\right) \tag{17}
\end{equation*}
$$

is valid, then the unperturbed motion is unstable.
Starting with the inequality $A>A_{0}+A_{1}+A_{2}$, we find that if

$$
\begin{equation*}
A_{2} \geqslant \frac{1}{2}\left(e^{-\alpha}-1\right)\left(1-e^{\beta}\right)-A_{1} \tag{18}
\end{equation*}
$$

then the unperturbed motion is unstable.
It is now obvious how one can obtain other sufficient conditions for stability or instability by the indicated method.

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